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APPLICATION OF WATERLOO MAPLE 9.5 AND WOLFRAM MATHEMATICA 5.1 SOFTWARE FOR ANALYTIC SOLVING OF CERTAIN NONLINEAR PARTIAL DIFFERENTIAL EQUATIONS OF PHYSICS

In the current paper some applications of the packet MAPLE (v. 9.5) for analytic solving of certain nonline partial differential equations have been presented. Additionally, for graphic presentation of the found solutions packet MATHEMATICA (v. 5.1) has been applied.

Keywords: nonlinear partial differential equations, analytic solutions, exact solutions, Navier-Stokes equation, Monge-Ampere equation, steady isentropic flow equation

1. Introduction

Nonlinear partial differential equations (NPDE’s) give a description of many interesting and important phenomena in physics. Some examples may be: Einstein equations

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in General Relativity or Navier-Stokes equations. There exist some powerful analytic
tools of solving such equations: Bäcklund transformation method [6], [4], symme-
try analysis method, [2], [4], a concept of strong necessary conditions [5], [6],
extended homogeneous balance method, [1], etc. However, it is possible they will not give all
solutions of investigated equation. The applied method, firstly presented in [7], gives
an opportunity for finding some wide classes of the solutions of certain NPDE’s. Pre-
sentation of examples of application of this method for some such equations, is the
main aim of this paper. This paper is organized as follows. In the next section we show
the equations, which we want to solve. In section 3 we present shortly the method.
Section 4 is devoted to show the results of applying of the mentioned method for
several NPDE’s. In section 5 some conclusions are included.

2. The equations

2.1. Euler and Navier-Stokes equation in dimensions (2+1)

Navier-Stokes equation has a form (for a plane flow), [4]:

\[(\Delta u)_t + u_y(\Delta u)_x - u_x(\Delta u)_y = n\Delta\Delta u, \quad u \in R\]  \hspace{1cm} (1)

where \(u \in R\) is a stream function and \(n\) is a viscosity coefficient. Euler equation
describes the flow of an ideal fluid, so in its case: \(n = 0\).

2.2. Homogeneous Monge-Ampere equation in the dimensions (3+1)

This equation has the following form, [2]:

\[\det(u_{\mu\nu}) = 0\]  \hspace{1cm} (2)

where \(u \in R, u_{\mu\nu} = \frac{\partial^2 u}{\partial x_\mu \partial x_\nu}, \mu, \nu = 0, 1, 2, 3\).

2.3. Equation of steady isentropic flow in the dimensions (2+0)

This equation has the form, [4]:

\[\left( g^2 - \left(\frac{\partial u}{\partial x}\right)^2 \right) \frac{\partial^2 u}{\partial x^2} - 2 \frac{\partial u}{\partial x} \frac{\partial u}{\partial y} \frac{\partial^2 u}{\partial x \partial y} + \left( g^2 - \left(\frac{\partial u}{\partial y}\right)^2 \right) \frac{\partial^2 u}{\partial y^2} = 0,\]  \hspace{1cm} (3)

where \(u \in R\) is a velocity potential and \(g = g(\left(\frac{\partial u}{\partial x}\right)^2 + \left(\frac{\partial u}{\partial y}\right)^2)\) is a sound speed. We
investigate the case, when \(g = \sqrt{\left(\frac{\partial u}{\partial x}\right)^2 + \left(\frac{\partial u}{\partial y}\right)^2}\).
3. An algorithm

3.1. Decomposition method

We want to find certain analytic solutions of some NPDE, [7]:

\[ F(x^\mu, u_1, \ldots, u_m, u_1 x^\nu, u_1 x^\alpha x^\nu, \ldots, u_m x^\mu, u_m x^\nu x^\mu, \ldots) = 0, \]
\[ \mu, \nu = 0, 1, 2, 3, \]

where \( m \in \mathbb{N} \) and \( u \) is in general an unknown function of class \( C^k, k \in \mathbb{N} \).

Let’s limit in this introduction to the case: \( m = 1, k = 2 \). We apply the following procedure, [7]:

I. we check, whether there is possible to make a decomposition of this equation into such fragments, which are characterized by a homogenity of the derivatives of the unknown function \( u \). “Homogenity of the derivatives”: these fragments should be products of at least two factors:

a. arbitrary expression, which may depend on the unknown function \( u \), its derivatives and the independent variables (in general: \( x^\mu, \mu = 0, 1, 2, 3 \)),

b. a sum of: at least two derivatives and (or) at least two products of the derivatives, so that there we can find for any derivative and (or) any product some other derivative and (or) other product, which has the same order and (or) degree.

II. For example, the investigated equation may be as follows:

\[ F_1 \cdot [(u_x)^2 + (u_y)^2] + F_2 \cdot [u_{xx} + u_{xy}] = 0, \]

where \( F_1 \) and \( F_2 \) may depend on \( x^\mu, u, u_{x^\nu}, \ldots \).

III. We insert some ansatzes, presented below. We collect all algebraic terms, appearing by all derivatives of \( f \) and require vanishing of them.

3.2. Ansatzes

1. The ansatz of first kind This ansatz has the form, [7]:

\[ u(x^\mu) = \beta_1 + f(a_\mu x^\mu + \beta_2, b_\nu x^\nu + \beta_3, c_\rho x^\rho + \beta_4), \]

where: \( v_\alpha x^\alpha = -v_0 x^0 + v_k x^k, k = 1, 2, 3 \). The parameters \( a_\mu, b_\mu, c_\mu, \mu = 0, 1, 2, 3 \), are some constants to be determined later, \( \beta_j, j = 1, \ldots, 4 \), are arbitrary constants and \( x^\mu \in R \). As well is in the case of combined ansatz.

2. The ansatz of second kind This ansatz is as it follows, [7]:

\[ u(x^\mu) = \beta_1 + f(A_0 x^0 x^1 + A_1 x^0 x^2 + A_2 x^0 x^3 + A_3 x^1 x^2 + A_4 x^1 x^3 + A_5 x^2 x^3 + \beta_2) \]

where \( A_\lambda, \lambda = 0, 1, \ldots, 5 \) are some constants to be determined later, \( \beta_j, x^\mu \) are as in the case of first ansatz.
3. The combined ansatz The form of this ansatz is, [7]:

\[
    u(x^\mu) = \beta_1 + f(a_{\mu}x^\mu + \beta_2, A_0x^0x^1 + A_1x^0x^2 + A_2x^0x^3 + A_3x^1x^2 + A_4x^1x^3 + A_5x^2x^3 + \beta_2),
\]

where \( a^\mu, \ldots, A^\lambda, \beta_j, \mu = 0, 1, 2, 3, \lambda = 0, \ldots, 5 \) and \( x^\mu \) are as in the cases of first and second ansatz.

Some important remarks:

1. These all ansatzes have a common feature: they contain some arbitrary function \( f \) of some appropriate arguments. It is the function of class \( C^2 \). We want to find wide class of the solutions as far as possible.

2. Of course, the number of the variables \( x^\mu \) in these ansatzes depends on the investigated equation.

3. The notation, concerned to the investigated equations, is: \( x^0 = t, x^1 = x, x^2 = y, x^3 = z \).

4. If the specification of the investigated equation requires that \( u \) has to be real function and from the further computations we obtain complex values of the coefficients \( a_{\mu}, \ldots, A_\lambda \), then we try to apply such ansatz:

\[
    u(x^\mu) = \beta_1 + \Re(f),
\]

where the dependence of the function \( f \) on the independent variables \( x^\mu \) is given by the ansatz.

### 3.3. Applied software and hardware

All computations and figures were done on the computers: SunFire 6800 (“saturn” in ACK-CYFRONET-AGH in Kraków), by using Waterloo MAPLE 9.5 Software and Wolfram MATHEMATICA 5.1 Software respectively, (grant No. MNiI/Sun6800/WSP/008/2005) and IBM Blade Center HS21 “mars”, by using Waterloo MAPLE 12 Software, (grant No. MNiSW/IBM BC HS21/AP/057/2008). All the mentioned computations were symbolic. The presented figures were made by applying a package “Graphics ’Animation’” and the function “MoviePlot3D” of MATHEMATICA 5.1 software.

### 4. Results of application of decomposition method by using MAPLE 9.5 and MATHEMATICA 5.1 software

#### 4.1. Solutions – existence of them and values of the coefficients of the ansatzes

As a result of inserting of the above ansatzes into each of the investigated equations, we obtain some system of algebraic equations, where the unknown quantities are the coefficients \( a_{\mu}, \ldots, A_\lambda \). We have got such system by using a MAPLE instruction
“collect”, thanks to it we can collect, according to the procedure, described in subsect. 3.1, all terms, which are by each derivative of the function $f$:

$$\text{collect(izentr, } [D[1](f), D[2](f), D[3](f), D[1,1](f), D[1,2](f), D[1,3](f),\
D[2,2](f), D[2,3](f), D[3,3](f)], \text{ distributed});$$

where “izentr” is a name of an Maple input, which contains an expression, created by mentioned inserting of the ansatizes. Moreover, there some short notation has been used: $D[i](f), D[i,j](f)$ denote $i$-th and $ij$-th derivative of the function $f$ respectively. Next we equal such collected expression to zero and we obtain an algebraic equation. For example we give here a fragment of such system, obtained in the case of the ansatz of first kind (6), applied for the equation of isentropic flow (3):

$$(a_1c_2 - a_2c_1)^2 = 0, \quad (9)$$
$$(b_1c_2 - b_2c_1)^2 = 0 \quad (10)$$

If we can write down and solve such system for the investigated equation, we know the exact solution, given by the ansatz, which we applied previously. For solving such system of algebraic equations we have applied the MAPLE instruction “solve”. The values of the coefficients for each of the investigated equations have been collected in the Table 1.

**Table 1**
The values of the coefficients. The coefficients, appearing at the right sides of the formulas, are arbitrary real constants. However, of course, if they are the denominators, they must be nonzero constants

<table>
<thead>
<tr>
<th>Equation</th>
<th>Solution of 1st kind</th>
<th>Solution of 2nd kind</th>
<th>Combined solution</th>
</tr>
</thead>
<tbody>
<tr>
<td>Euler- and Navier-Stokes</td>
<td>unphysical solution: $a_2 = ia_1, b_2 = ib_1, c_2 = ic_1$ (see subsect. 3.1 and some comments beneath)</td>
<td>the solution does not exist (see subsect. 3.1 and some comments beneath)</td>
<td>the solution does not exist (see subsect. 3.1 and some comments beneath)</td>
</tr>
<tr>
<td>Monge-Ampere</td>
<td>$a_\mu, b_\mu, c_\mu \in R$ - arbitrary</td>
<td>$A_0 \in R$ - arbitrary, $A_1 = \frac{A_2A_3}{A_4}, A_5 = 0$</td>
<td>$a_3 \in R$ - arbitrary, $A_1 = \frac{a_2A_0}{a_1}, A_4 = \frac{a_3A_5}{a_2}, A_3 = 0$</td>
</tr>
<tr>
<td>isentropic flow</td>
<td>$a_1 = \frac{a_2b_1}{b_2}, c_1 = \frac{b_1c_2}{b_2}$</td>
<td>solution does not exist</td>
<td>solution does not exist</td>
</tr>
</tbody>
</table>

Several comments on the solutions of Euler- and Navier-Stokes equations (1).

1. It is clear that all real harmonic functions are the solutions of Euler- and Navier-Stokes equations (1), [4]. From the viewpoint of the method presented in this paper, it is the special case of this method. Further, the parameters $a_k, b_k, c_k, k = 1, 2$ collected in the Table 1, cause the solution is complex. However, although this fact and the non-existence of the solutions of second kind and combined
solutions for Euler- and Navier-Stokes equations, we can try to apply the point 4 of some important remarks from the previous section. Hence, we can show that there in this case exists more general (and real) solution, than the original solution of first kind, (some suggestion is included in [4]):

$$u(x, y, t) = \beta_1 + f(x^0, a_k x^k + \beta_2, b_k x^k + \beta_3, c_k x^k + \beta_4) +$$

$$+ f(x^0, a_k^* x^k + \beta_2^*, b_k^* x^k + \beta_3^*, c_k^* x^k + \beta_4^*)$$  \hspace{1cm} (11)

where $a_k, b_k, c_k, k = 1, 2$ are given in Table 1 and “*” denotes complex conjugation.

2. Besides there exists semi-second kind solution:

$$u(x, y, t) = \beta_1 + h_1(x^0)x^1 + h_2(x^0)x^2 +$$

$$+ h_3(x^0)x^1 x^2 + h_4(x^0) \cdot ((x^1)^2 - (x^2)^2)$$  \hspace{1cm} (12)

where $h_j(x^0) \in R, j = 1, 2, 3, 4$ are arbitrary functions of class $C^2$. This is such, because there in Laplace equation is no differentiation with respect to the variable $x^0$ and thanks to this fact we can put here arbitrary functions of this variable.

3. The most general solution, which we can give here for Euler- and Navier-Stokes equations (1), from the viewpoint of the ansatz of second kind and combined ansatz, has the form:

$$u(x, y, t) = \beta_1 + h_1(x^0)x^1 + h_2(x^0)x^2 + h_3(x^0)x^1 x^2 +$$

$$+ h_4(x^0) \cdot ((x^1)^2 - (x^2)^2) + f(x^0, a_k x^k + \beta_2, b_k x^k + \beta_3, c_k x^k + \beta_4) +$$

$$+ f(x^0, a_k^* x^k + \beta_2^*, b_k^* x^k + \beta_3^*, c_k^* x^k + \beta_4^*),$$  \hspace{1cm} (13)

where $a_k, b_k, c_k, k = 1, 2$ are given in Table 1 and $h_j(x^0) \in R, j = 1, 2, 3, 4$ are arbitrary functions of class $C^2$. As we see, in this case we have a superposition of (11) and (12).

4.2. Figures of the semi-combined solution of Navier-Stokes equation

Hereunder we present the figures of a fragment of an animation of the semi-combined solution (13), when $h_1(t) = h_2(t) = h_3(t) = 0, h_4(t) = \exp(-t^2)$, $a_1 = 2, b_1 = 4, c_1 = 5, \beta_i = 0, i = 1, 2, 3, 4, x \in (-5, 5), y \in (-10, 10), t \in (10, 1000)$ and the number of plot points was equal to 50.

$$u(x, y) = \exp(-t^2) \ast (x^2 - y^2) +$$

$$+ \exp(\coth(\tan(t) + (2x + 2iy) + (4x + 4iy)(5x + 5iy))) +$$

$$+ \exp(\coth(\tan(t) + (2x - 2iy) + (4x - 4iy)(5x - 5iy)))$$  \hspace{1cm} (14)

Because of limited space of this paper, we present only four figures of the animation.

In the order to do this animation we used the package “Graphics ‘Animation’” and the function “MoviePlot3D” of MATHEMATICA.
Three first figures (Figs 1 and 2) present an evolution of the mentioned solution within the interval \( t \in (10, 1000) \) and on the four figure we see the solution at the end of the interval, when \( t = 1000 \).

In [3] so called \( 2\delta \) pulse solution of double sinus-Gordon equation was presented. The first three figures show something, which is very similar to a collision of such two \( 2\delta \) pulse solutions, however, of course it is to be investigated later, [8]. If we look at the Figure 3, which present the solutions at the end of time interval, we can say that the presented solutions seem to not decay in time, but obviously, a question about their stability is open [8].

![Fig. 1. Solution (14)](image)

5. Conclusions and future work

As far as the solution of first kind of Monge-Ampere equation (2) is concerned: this solution has been presented in [2], but by applying the symmetry analysis method.

The presented method gives for the investigated equations enough wide classes of exact solutions and it seems to be easier for applications, than for example symmetry analysis method. Moreover, if we look on the ansatzes (6)–(8) may be generalized to the ansatzes dependent on more arguments, like these ones, occuring in (6)–(8).

From the other hand, this method requires some carefulness, when there is required that the solution has to be real function, because often the roots of the system of algebraic equations turn to be complex numbers and sometimes it is easy to obtain unphysical solutions. This paper is concerned only several examples of application of the decomposition method. Future work will be related to a generalization of the ansatzes, extension of this method for other families of NPDE’s and to investigation of the physical features of obtained solutions, [8].
Fig. 2. Solution (14)

Fig. 3. Solution (14) in $t = 1000$
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[8] Ł. T. Stępień, work in progress